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Research Paper

Turbulent Forced Convection and Entropy Generation of Impinging Jets of Water-Al₂O₃ Nanofluid on Heated Blocks

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Abstract. A computational analysis on water-Al₂O₃ nanofluid turbulent forced convection is performed to analyze heat transfer and entropy production in a channel containing heated blocks, cooled by impinging jets. The two phase mixture model (TPMM) is used. This work can be considered as a good contribution to improving thermal performance and enhancing the heat transfer rate (HTR) for some engineering industries. After completing the current study, we concluded that the increase in the Reynolds number (Re) and the volume fraction of nanoparticles (φ), the decrease in spacing between the heated block (D_b) and moving the location of the second jet (J2) to the first jet (J1) contribute to increasing the heat transfer rate (HTR). In addition, the TPMM gives higher values of average Nusselt number \overline{Nu} than the single-phase model (SPM). The thermal ($\dot{S}_{g,th}$), frictional ($\dot{S}_{g,v}$) and total ($\dot{S}_{g,t}$) entropy generation values increase with Re and φ . When D_b is reduced, $\dot{S}_{g,t}$ increases. However, $\dot{S}_{g,t}$ increases when the jet position vary from J2 to J1. Different correlations are proposed for \overline{Nu} . Our results are compared with data available in the literature.

Keywords: Heat transfer, Turbulent flow, entropy generation, nanofluid, impinging jets, heated blocks.

1. Introduction

Increased HTR is an active and important area of engineering research due to its contribution to developing highly efficient cooling technologies to meet ever-growing demand for developing efficient designs for thermal equipment. Various techniques have been proposed in recent years to enhance convective heat transfer [1-4]. In order to improve the HTR in different systems, nanofluid have been used more widely in recent years compared to conventional heat transfer fluids [5-6]. In this regard, one of the most promising strategies for promoting heat transfer compared to conventional forced convection techniques is its liquid jet effect cooling due to its performance the coefficient of heat transfer (h) is higher. As an obvious implication, liquid jets were used in various industrial and engineering applications for example cooling of heat engines, thermal control in high-heat-dissipation electronic devices and thermal treatment of metals and materials processing [7]. Several experimental and numerical studies have been conducted, for example. Menni et al. [8] studied the effect of different base fluid on the HTR. They found that the pure ethylene glycol flows give the largest thermal exchange compared with the water and water- ethylene glycol mixture. Ibrahim and Ahmed [9] investigated numerically the impinging jet's laminar heat transfer with nanofluid. The outcomes demonstrate that the $\overline{\text{Nu}}$, performance factor, $\dot{S}_{q,t}$ increases with increasing φ . Menni et al. [10,11] studied the effect the presence of nanoparticles and baffle plates in a channel, they also examined the effect of different outlet models in baffled channels on the HTR used Al₂O₃-H₂O. A detailed review of the numerical and experimental studies performed to improve the HTR by Menni et al. [12-13]. Abdelrehim et al. [14] studied SPM and TPMM of nanofluid on heat transfer properties. Results show that the TPMM gives higher values of Nu and Nu . Sheikholeslami et al. [15-16]. Used TPMM to simulate fluid flow, heat transfer, energy and entropy evaluation inside a tube with enhanced turbulator. Selimefendigil and Öztop [17] investigated numerically the effects of different parameters such as pulse frequency, Re and φ . They showed that in the pulsating flow case, HTR will improve with increasing Re



and φ . A numerical study using nanofluids to cool an isothermal hot surface with an adiabatic rotating cylinder was carried out by Selimefendigil and Öztop [18]. It was found that the Nu increases with Re, but the HTR was higher in the case when the cylinder is closer to the jet inlet. The influences of uniform and non-uniform impact jets velocity were carried out by Izadi et al. [19]. Results indicate that raising the values of φ enhances HTR. Additionally, when using non-uniform impingement jet with a low speed distribution, the thermal performance is improved. Lamraoui et al. [20] investigated numerical the flow characteristics of a non-Newtonian Al₂O₃-water nanofluid. It was found that Nu is much higher for non-Newtonian nanofluid than Newtonian flow. Manca et al. [21] analyzed fluid dynamic and thermal behaviors of an enclosed laminar impinging slot jets with nanofluid. Results show that h and Nu values increase with increasing φ and Re. Amjadian et al. [22] experimentally studied the properties of the impact jet on a hot surface using Cu₂O-water nanofluid. Results showed that the rise in Re and φ increases the h value. Yousefi et al. [23] carried out an experimental study into the impingement of a planar jet on a plate. They showed a decrease in h with an increase in Re from 1732 to 2250, also they found that the use of Al₂O₃ - nanofluid at low φ of 0.02% and 0.05% improved h and \bar{h} . Teamah et el. [24] studied numerically and experimentally the behavior of nanofluids flow impingement on horizontal flat plate. They found that increasing φ increases h compared to pure water. The heat transfer properties of air/nanofluid jet cooling flux on a rotating hot circular disk, using a multiphase VOF model were studied by Mahdavi et al. [25]. Results showed that Nu values increase with increasing φ and disk rotation speed. Sun et al. [26] experimentally studied heat transfer with Cu Nanofluids from single impinging jet. It was found that the coefficient h is higher when using round nozzles compared to square nozzles. Pratap et al. [27] conducted the effect of nanofluids for the enhancement of the HTR by impingement jet. They observed that the Nu values increases with the increase in H/D ratio. A numerical study of the impact jet in a small channel using the Eulerian-Eulerian biphasic method was performed by Venkatasubbaiah and Abhijith [28]. HTR enhancement is 25% for Al₂O₃ water nanofluid and 60% for Cu-water nanofluid compared to pure water, at Re = 300 and $\varphi = 5\%$. Abanti et al. [29] studied the effects of heat transfer properties on a moving plate. They found that the \overline{Nu} significantly rises with Re and with the plate velocity. A numerical simulation in a confined jet was performed by Rahimi-Esbo et al. [30]. Increasing Re, the aspect ratio and φ lead to an increase in the values of $\overline{\,{
m Nu}}$. Results showed that raising the values of Re, aspect ratio and $\,arphi$ conducts an increase in values of $\overline{\text{Nu}}$. Youssefi Lavouraki et al. [31] conducted a numerical simulation of confined jet using TPM. The outcomes indicate that Cf and Nu decrease and increase, respectively, with a higher of Re, H/W ratio and φ . Youssef Lavouraki et al. [32] investigated the heat transfer on a confined jet. Results indicated that an increase in φ contributes to an increase in Nu. Paulraj and Sahu [33] investigated numerically the conjugated heat transfer in laminar impinging jet on various hot sources using different nanofluids. They observed that the conjugate HTR increases with the decrease of nanoparticles diameter. It was also found that Al₂O₃-water nanofluid appears to be higher Nu compared to other nanofluids. The fluid flow and heat transfer properties resulting from the cooling of a group of heated blocks in a channel using an air jet was analyzed by Arquis et al. [34]. HTR of the small heated blocks is improved. Lam and Prakash [35] numerically studied heat transfer and \dot{S}_q of heat sources with or without porous media. They showed that $\overline{\text{Nu}}$, $\dot{S}_{g,t}$ and $\dot{S}_{g,v}$ increases with increasing Darcy number, Re and porous layer thickness. Paulraj et al. [36] performed relevant research into the flow and heat transfer of several slot jets affecting different heat sources. Results proved that HTR decreases when the channel height increases. Boudraa et Bessaïh [37] investigated the behavior of fluid flow and heat transfer around a hot cubic block exposed to a cross-flow and an impinging jet. It was observed that changing the jet's location toward the channel inlet could enhance the HTR. Nimmagadda et al. [38] studied the effect of uniform/non-uniform magnetic field. It was found that increasing the strength of the magnetic field improves the HTR.

According to our knowledge of numerical and experimental researches conducted in this area, this is the first study relevant to turbulent forced convection analysis of water-Al₂O₃ nanofluid flow, heat transfer and entropy generation in a channel containing an array of heated blocks mounted on the bottom wall and influenced by the TPMM impinging jet. The main objective concerns the effects of various parameters such as Re, φ , D_b and the optimal position of J2 on the HTR. Some correlations of $\overline{\text{Nu}}$ are obtained. This work can be improved thermal performance of some engineering industries.

2. Geometry description and mathematical model

The computational domain is a two-dimensional (2D) symmetric channel containing an array of heated blocks placed on the bottom wall and cooled by impinging jets (Figure 1). The geometric parameters are: L = 0.04m, H = L, $D_b = 3L/4$, h = L/4, N = L/2, L1 = 16L where L, h are the length and height blocks, D_b is the space between each two blocks. The jets enter through the nozzle of width N, where the distance between it (nozzle) and the bottom wall is H.

Table 1. Physical properties of the water and Al₂O₃ at T= 293 K.

Properties	Water	Al_2O_3
ρ	998.2	880
C_p	4182	36
μ	0.001003	3890
k	0.6	/



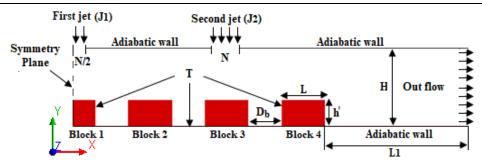


Fig. 1. Geometry.

The main properties of water and nanoparticles are given in Table 1 [17].

2.1 Assumptions

The heat transfer behavior of nanofluids is modeled using TPMM. Simplifying the study is made with the following assumptions:

- The flow is considered turbulent, incompressible and assumed to be a Newtonian fluid.
- The base fluid and nanoparticles to be in thermal equilibrium.
- The nanoparticles are assumed to be spherical shape, uniform in size and shape.
- The speed of the nanoparticles has the same speed as the base fluid.
- The dissipation of viscosity and approximation of Boussinesq are negligible.

2.2 Governing equations

Based on the assumptions mentioned above, the equations which govern the TPMM can be expressed as follows:

Continuity equation [39]:

$$\nabla \cdot (\rho_{\scriptscriptstyle m} V_{\scriptscriptstyle m}) = 0 \tag{1}$$

Momentum equation [16]:

$$\nabla \cdot (\rho_{\scriptscriptstyle m} \mathbf{V}_{\scriptscriptstyle m} \mathbf{V}_{\scriptscriptstyle m}) = -\nabla \mathbf{P} + \nabla \cdot (\nabla \mathbf{V}_{\scriptscriptstyle m} + \nabla \mathbf{V}_{\scriptscriptstyle m}^{\scriptscriptstyle \mathsf{T}}) \mu_{\scriptscriptstyle m} + \nabla \cdot \sum_{k=1}^{n} (\rho_{\scriptscriptstyle k} \varphi_{\scriptscriptstyle k} \mathbf{V}_{\scriptscriptstyle dr,k} \mathbf{V}_{\scriptscriptstyle dr,k})$$
 (2)

Energy equation [41]:

$$\nabla \cdot \sum_{k=1}^{n} (\rho_k \varphi_k \mathbf{V}_k \mathbf{C}_{p,k} \mathbf{T}) = \nabla \cdot (\mathbf{K}_m \nabla \mathbf{T}) \tag{3}$$

Volume fraction equation [42]:

$$\nabla \cdot (\varphi_{np} \rho_{np} \mathbf{V}_m) = -\nabla \cdot (\varphi_{np} \rho_{np} \mathbf{V}_{dr,np}) \tag{4}$$

The mixture velocity of nanofluid can be expressed as:

$$V_{m} = \sum_{k=1}^{n} \frac{\varphi_{k} \rho_{k} V_{k}}{\rho_{m}} \tag{5}$$

The drift velocity of the nanoparticle is:

$$V_{dr,k} = V_k - V_m \tag{6}$$

The slip velocity is:

$$V_{nv,bf} = V_{nv} - V_{bf} \tag{7}$$

The drift velocity and slip velocity can be related as:

$$V_{dr,np} = V_{np,bf} - \sum_{k=1}^{n} \frac{\varphi_k \rho_k V_{fk}}{\rho_m}$$
(8)

Following equations proposed by [43] and [44], the slip velocity ($V_{np,bf}$) and drag function (f_{drag}) are, respectively.

$$V_{w, M} = \frac{\rho_{w} d_{w}^{2}}{18 \mu_{w} f_{dwa}} \frac{(\rho_{w} - \rho_{w})}{\rho_{w}} g - (V_{w}.\nabla) V_{w}$$
(9)

$$f_{asy} = \begin{cases} 1 + 0.15 \, \text{Re}_{p}^{\text{olar}}, & \text{Re}_{p} \le 1000 \\ 0.0183 \, \text{Re}_{p}, & \text{Re}_{p} \ge 1000 \end{cases}, \, \, \text{Re}_{p} = \frac{V_{\text{m}} d_{p}}{\nu_{\text{m}}}$$
(10)



where V_a and d_a are the kinematics viscosity of mixture and nanoparticles diameter, respectively.

2.3 Turbulence model

For the standard $k-\varepsilon$ turbulence model, the turbulent kinetic energy (k) and the turbulent dissipation rate (ε) equations [45-46] are:

$$\nabla \cdot (\rho_{\scriptscriptstyle m} V_{\scriptscriptstyle m} k) = \left(\frac{\mu_{\scriptscriptstyle t,m}}{\sigma_{\scriptscriptstyle k}} \nabla k\right) + G_{\scriptscriptstyle k,m} - \rho_{\scriptscriptstyle m} \varepsilon \tag{11}$$

$$\nabla \cdot (\rho_{_{m}} V_{_{m}} \varepsilon) = \left(\frac{\mu_{_{t,m}}}{\sigma_{_{k}}} \nabla \varepsilon\right) + \frac{\varepsilon}{k} (C_{_{1}} G_{_{k,m}} - C_{_{2}} \rho_{_{m}} \varepsilon) \tag{12}$$

where $\,\mu_{\scriptscriptstyle {\rm t},m}\,$ and $\,{\rm G}_{\scriptscriptstyle {\rm k},m}\,$ are given as follows:

$$\mu_{\scriptscriptstyle t,m} = \rho_{\scriptscriptstyle m} C_{\scriptscriptstyle \mu} \frac{k^{\scriptscriptstyle 2}}{\varepsilon} \tag{13}$$

$$G_{k,m} = \mu_{t,m} \left(\nabla V_m + (\nabla V_m)^T \right) \tag{14}$$

 C_1 , C_2 , C_μ , σ_k and σ_ϵ are empirical constants: $C_1 = 1.44$, $C_2 = 1.92$, $C_\mu = 0.09$, $\sigma_k = 1.0$ and $\sigma_\epsilon = 1.3$.

The nanofluid properties are calculated as:

• Density [47]:

$$\rho_{\rm m} = (1 - \varphi)\rho_{\rm bf} + \varphi\rho_{\rm m} \tag{15}$$

• Specific heat [48]:

$$Cp_{m} = \frac{(1 - \varphi)Cp_{bf}\rho_{bf} + \varphi Cp_{np}\rho_{np}}{\rho_{m}}$$
(16)

• Dynamic viscosity [49]:

$$\mu_{\rm m} = (1 + 7.3\varphi + 123\varphi^2)\mu_{\rm bf} \tag{17}$$

• Thermal conductivity [50]:

$$\mathbf{k}_{m} = \mathbf{k}_{bf} \left[\frac{\mathbf{k}_{np} + 2\mathbf{k}_{bf} + 2\varphi(\mathbf{k}_{np} - \mathbf{k}_{bf})}{\mathbf{k} + 2\mathbf{k}_{bf} - 2\varphi(\mathbf{k}_{nm} - \mathbf{k}_{bf})} \right]$$
(18)

The total entropy generation includes the fluid friction and heat transfer terms are defined as [51]:

$$\dot{S} = \dot{S}_{a,th} - \dot{S}_{a,v} \tag{19}$$

$$\dot{S}_{g,th} = \frac{k_m}{T^2} \left[\left(\frac{\partial T}{\partial x} \right)^2 + \left(\frac{\partial T}{\partial y} \right)^2 \right]$$
 (20)

$$\dot{S}_{g,v} = \frac{\mu_m}{T} \left[2 \left[\left(\frac{\partial u}{\partial x} \right)^2 + \left(\frac{\partial v}{\partial y} \right)^2 \right] + \left[\left(\frac{\partial u}{\partial y} \right) + \left(\frac{\partial v}{\partial x} \right) \right]^2 \right]$$
(21)

Table 2. Boundary conditions.

Туре	U _i (m / s)	V _j (m / s)	T (K)
First jet inlet (J1)	0	$\mathbf{v}_{_{j}}$	293
Second jet inlet (J2)	0	v_{j}	293
Outflow	$\frac{\partial u}{\partial x} = 0$	$\frac{\partial \upsilon}{\partial x} = 0$	$\frac{\partial T}{\partial x} = 0$
Symmetric side	0	$\frac{\partial \upsilon}{\partial x} = 0$	$\frac{\partial T}{\partial x} = 0$
wall	0	0	$\frac{\partial T}{\partial x} = \frac{\partial T}{\partial y} = 0$
Temperature	0	0	343



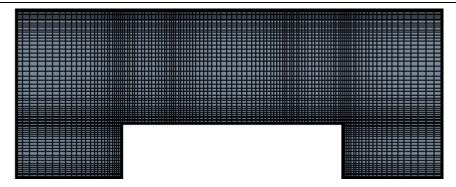


Fig. 2. Computational grid.

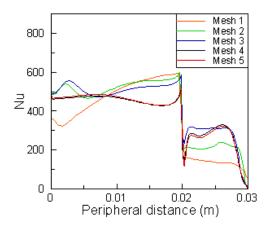


Fig. 3. Variation of local Nu for the first block with different tested grids at $\varphi = 5\%$ and Re = 20000.

2.4 Boundary conditions

The boundary conditions used are summarized in Table 4. At the jets inlet (J1 and J2), it was assumed that the flow is turbulent at a uniform velocity (V_i). At the channel exit, all gradients are equal to zero, the upper and bottom wall except for the hot part are adiabatic with nonslip condition, the distance around the blocks and between them are assumed under constant temperature with nonslip condition. At the symmetric side, we assumed that zero normal gradients of all variables.

3. Computational model

The Ansys-Fluent 14.5 software [52] was used for simulating the turbulent heat transfer flow. We have used second upwind order scheme and the SIMPLE algorithm. In order to confirm the negligible effect of the grid on the results performed by raising the mesh density successively until further refinement revealed a difference between the consecutive results of less than 1% and according to the best compromise between precision and available calculation means, a grid of approximately 110.000 structured quadrilateral cells (1100×100 in the x and y directions, respectively) was adopted for all cases we studied as shown in Figure 2. This mesh was designed to capture the sharp gradients and boundary layers near the wall and the blocks.

3.1 Grid independence study

The mesh independence test was performed for Al2O3–water nanofluid with Re = 20000 and φ = 5% . Figure 2 shows the distribution of Nu along the first heated block calculated on five (5) different meshes (see table 3). After comparing the five cases, the mesh number 4 was chosen as an acceptable mesh.

4. Results and discussion

A numerical analysis related to the turbulent forced convection of water-Al₂O₃ nanofluid is carried out to analyze the fluid flow, heat transfer and entropy generation in a channel containing heated blocks, mounted on the bottom wall and cooled by impinging jets. The parameters are Re, φ , D_b and the position optimal of the J2. We present Nu, $\overline{\text{Nu}}$, $\dot{S}_{g,th}$, $\overline{\text{Cf}}$, $\dot{S}_{g,t}$ and Be. Various correlations for $\overline{\text{Nu}}$ are presented.

Table 3. Number of nodes for various meshes.

Meshes	1	2	3	4	5
Number of nodes	715x70	875x80	1000x90	1000x100	1217x115



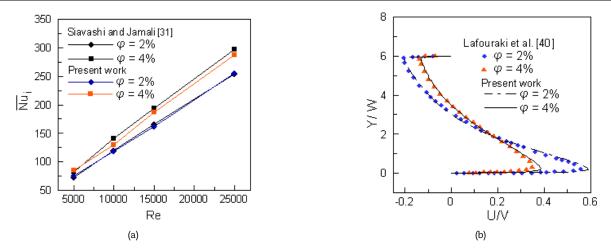


Fig. 4. Comparison of present research with data available.

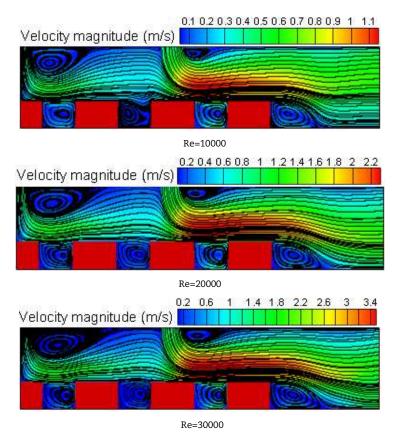


Fig. 5. Effect of Re on stream function and velocities magnitude contours for different Re at $\varphi=5\%$.

4.1 Validation

In order to confirm the validity of the fluid flow and heat transfer solution, we used the standard $k-\varepsilon$ turbulence model with enhanced wall treatment by employing TPMM. Our results are compared with [31] and [40], and are shown in Figure 4a and Figure 4b, respectively. The first figure (4a) shows a good agreement with the work of Siavashi and Jamali [31] by comparing the $\overline{\text{Nu}}$ at different Re (Re = 5000 to 25000). With regard the second figure (4b), a similar trend is observed with the results of Lafouraki et al. [40], based from the comparison of the velocity magnitude at $\varphi = 2\%$ and $\varphi = 4\%$.

4.2 Effect of Reynolds number

Figure 5 represent the streamline and magnitude velocities contours one on the other at different Re (Re =10000, 20000 and 30000) using Al₂O₃-water nanofluid with $\varphi=5\%$. It can be noted that the size of the vortex formed at neighborhood jets increases with Re, where the J1 is located in the middle of the upper surface of the symmetric channel and J2 located on the middle of the third block. The increase in Re results in an increase of the intensity of the vortices formed in the channel between the heated blocks, after the last block the recirculation zone formed due to sudden expansion, this vortex also increases with Re. A higher streamline concentration (where speed takes great values) was found at the top of the blocks 3 and 4 due to the dominance of the J2.



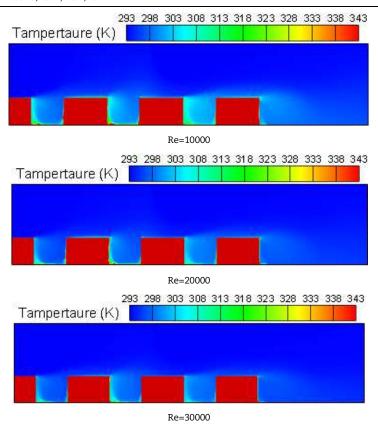


Fig. 6. Effect of Re on isothermal contours, at $\varphi=5\%$.

Figure 6 presents the effect of Re (Re=10000, 20000 and 30000) on the temperature distribution around the heated blocks, at $\varphi=5\%$. We can observe that the isotherms occupy a considerable part of channel around and between the heated block for a low Re (Re=10000), because of that momentum of nanofluid not enough to decrement thermal layer. Note the opposite when increasing Re. This means that the thermal boundary layer thickness gradually decreases when Re increases.

Figure 7 illustrates the distribution of Nu in terms of Re around four (4) heated blocks, at $\varphi = 5\%$. Results show that an increase in Re allows an increase in Nu at the top of each block due to a reduction in thermal boundary layer thickness. Figure 7 (a-d) shows the increase of Nu on the left and right sides of the blocks with the increase of Re. This is due to the increase in the intensity of the vortices that form between each block in the channel. We observe that the Nu distribution around block 2 is the same as block 4 with a difference in values of Nu. The maximum values of Nu are located around the last block due to deflection of the J2 and the domination of this jet (J2) compared to J1.

Figure 8 depicts the variation of $\overline{\text{Nu}}$ and $\overline{\text{Cf}}$ around each heated block, at $\varphi = 5\%$ and different Re. It can be noted that the increase in Re is accompanied by an increase $\overline{\text{Nu}}$ and a reduction in $\overline{\text{Cf}}$. Decrease in $\overline{\text{Cf}}$ at high Re is due to decreases in boundary layer thickness, so the velocity and temperature gradients increase. This leads to increase in $\overline{\text{Cf}}$. This means that the HTR is improved.

4.3 Effect of volume fraction

Figure 9 shows the variation of average $\overline{\text{Nu}}$ around of each heated block with different φ , at Re=20000. We can observe that all nanofluids have higher $\overline{\text{Nu}}$ than those of the pure water ($\varphi=0$), because as we know the increase in φ improves thermophysical properties of the Al₂O₃-water nanofluid. Also, we can note that the $\overline{\text{Nu}}$ increases with increasing φ , because of the large exchange of energy due to the chaotic motion of nanoparticles, hence HTR is enhanced with increase in φ . We can conclude that the maximum values of $\overline{\text{Nu}}$, located around the last block, is due the domination of the J2 compared to J1.

4.4 Effect of the second jet position

To choose the optimal position of J2 to improve the HTR, we changed the position of J2 several times (see Table 4.) along the upper wall. Figure 10 illustrates the variation of $\overline{\text{Nu}}$ for different positions of the J2, at $\varphi=3\%$ and different Re. We can note that HTR increases with increasing in Re due to the increase of the length and strength of the recirculation zones formed between the heated blocks, which allows decreasing the thermal boundary layer thickness, so HTR is improved. Changing the position of the J2 towards the J1 also enhances HTR. This allows us to take advantage of the lost momentum of the J2, which increases the temperature and velocity gradients along the heated blocks. We can conclude that the optimal location for the J2, which gives the highest HTR is located on the second block.



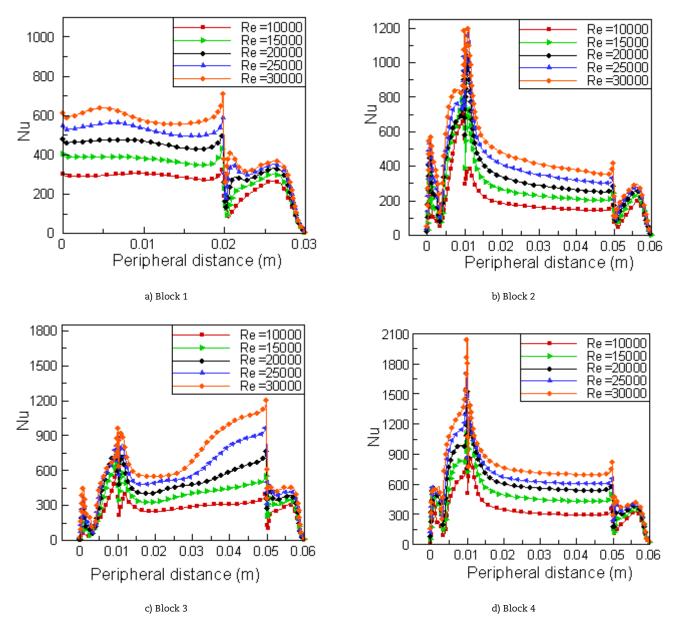


Fig. 7. Variation of Nu with peripheral distance and Re for different blocks, at $\varphi = 5\%$.

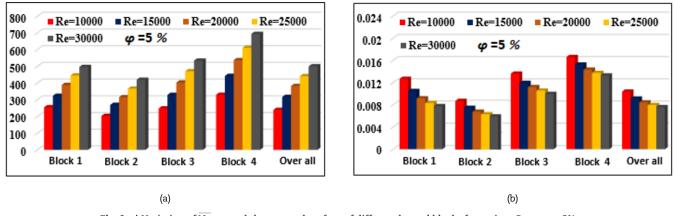
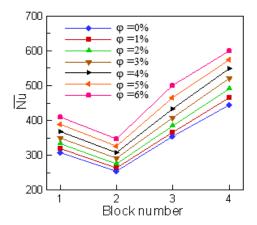


Fig. 8. a) Variation of $\overline{\text{Nu}}$ around the exposed surface of different heated blocks for various Re, at $\varphi=5\%$ b) Variation of $\overline{\text{Cf}}$ around the exposed surface of different heated blocks for various Re, at $\varphi=5\%$





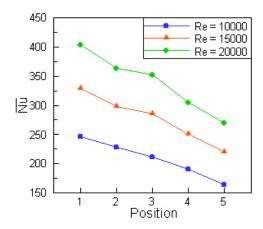
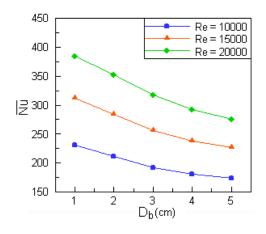


Fig. 9. Variation of $\overline{\text{Nu}}$ around the exposed surface of different heated bloc ks for various φ , at Re=20000.

Fig. 10. Variation of $\overline{\text{Nu}}$ along the heated blocks for different position of J2 at $\varphi=3\%$ and different Re.



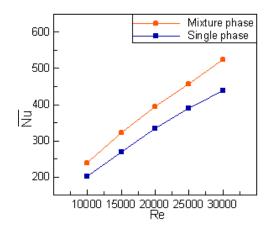
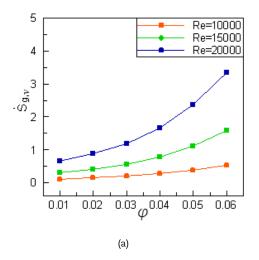


Fig. 11. Variation of $\overline{\text{Nu}}$ along the heated blocks for different D_{b,} at $\varphi=$ 5% and different Re.

Fig. 12. Comparison of $\overline{\text{Nu}}$ for different heated blocks between mixture and single phase model, at different Re and $\varphi=5\%$.



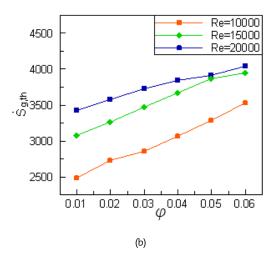


Fig. 13. The effect of φ on a) $\dot{S}_{g,th}$ and b) $\dot{S}_{g,v}$

Table 4. Different distances between J1 and J2.

Position	1	2	3	4	5
Distance between J1 and J2 (cm)	7.5	11.5	14	17.5	21



4.5 Effect of spacing between heated blocks

Figure 11 depicts the effect of D_b on \overline{Nu} of Al_2O_3 —water nanofluid, at $\varphi=5\%$ with different Re. The position of the second jet is fixed on the third block despite the change of D_b . We found that the increase in Re contributes significantly to the improvement in the HTR. So the influence of the jets on the heated blocks leads to increases the temperature and velocity gradients, when we increase in D_b we get the opposite, which means the HTR decreases.

4.6 Comparison between single and two-phase model

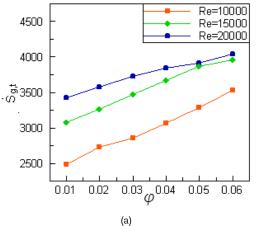
Figure 12 presents the comparison of the $\overline{\text{Nu}}$ for different Re at $\varphi = 5\%$ with both SPM and TPMM. According to these results, we can observe that the values of $\overline{\text{Nu}}$ increase with the increase of Re, where the values of $\overline{\text{Nu}}$ in the TPMM model are higher than the SPM, the same results were found by [14, 53].

4.7 Entropy generation analysis

Figure 13 illustrates the φ effect on $\dot{S}_{g,th}$ and $\dot{S}_{g,v}$ at different Re. Results revealed that the $\dot{S}_{g,th}$ and $\dot{S}_{g,v}$ values increase with the increase of φ due the improvement of thermal conductivity and viscosity of fluid by adding nanoparticles. The values of both $\dot{S}_{g,th}$ and $\dot{S}_{g,v}$ are also increasing due the increased in velocity and temperature gradients with the increase in Re. So, $\dot{S}_{g,th}$ and $\dot{S}_{g,v}$ increasing with the increase in φ and Re.

Figure 14a illustrates the φ effect on $\dot{S}_{g,t}$ at different Re. Results show that the values of $\dot{S}_{g,t}$ increase with the increase in Re and φ , with regard to Figure 11b, which illustrates the φ effect on Be at different Re. The increase in φ and Re leads to reduce the values of Be, this decreases is due to the increase in $\dot{S}_{g,v}$ and $\dot{S}_{g,th}$ with Re and φ (See Fig. 11a and Fig. 11b). From Figures 11a and 10b, and by using the definition of Be (Be = $\dot{S}_{g,th}$ / $\dot{S}_{g,t}$), we can explain that lower values obtained by saying that the $\dot{S}_{g,th}$ values are much greater than $\dot{S}_{g,v}$ ($\dot{S}_{g,th}$ dominate on $\dot{S}_{g,v}$). In [54], the same results were obtained.

Figure 15 shows the variation of $\dot{S}_{g,t}$ for different position of J2, at $\varphi=3\%$ with different Re. We can note that $\dot{S}_{g,t}$ increases with increasing in Re due the increased in velocity and temperature gradients with the increase in the values of Re. Also, we can clearly notice that by changing the position of the second towards the first jet, this allows us to take advantage of the lost momentum of the J2, which increases the temperature and velocity gradients along the heated blocks. So the $\dot{S}_{g,t}$ increases and takes its highest values over the second block.



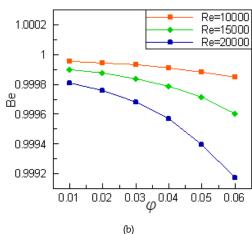


Fig. 14. The effect of φ on a) \dot{S}_{a+} and b) Be

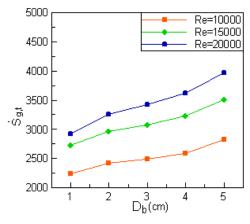


Fig. 15. The effect of D_b on \dot{S}_{at} .



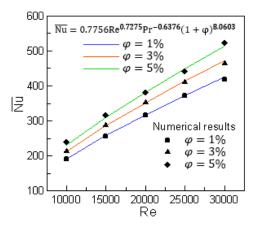


Fig. 16. Correlation of $\overline{\text{Nu}}$ for different Re and φ .

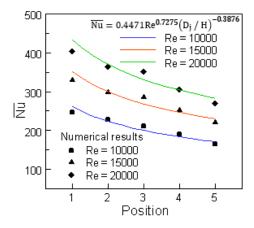


Fig. 18. Correlation of the \overline{Nu} for different Re and D_j .

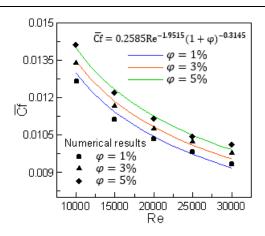


Fig. 17. Correlation of $\overline{\mathsf{Cf}}$ for different Re and φ .

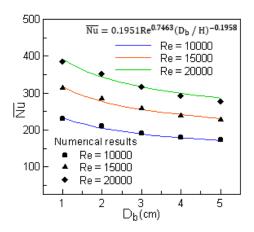


Fig. 19. Correlation of Nu for different Re and Db.

4.8 Correlations

The first correlation proposed is related of the $\overline{\text{Nu}}$ along the heated blocks, presented as follows:

$$\overline{\text{Nu}} = 0.7756 \,\text{Re}^{0.7275} \,\text{Pr}^{-0.6376} (1+\varphi)^{8.0603}$$
 (22)

and it is validated with the numerical results for an error between -1.5% and 3.5%, as shown in Figure 16. This correlation is defined as function of the Re, Pr and φ ($CRe^n Pr^m (1+\varphi)^k$), where $10000 \le Re \le 30000$ and $1\% \le \varphi \le 5\%$. In addition, this correlation can help engineers working in the industry who need to design this configuration.

The second correlation that we proposed relates to \overline{Cf} along the heated blocks, presented as follows:

$$\overline{\text{Cf}} = 0.2585 \,\text{Re}^{-1.9515} (1 + \varphi)^{-0.3145}$$
 (23)

and it is available with the numerical results for an error not exceeding 3%, as shown in Figure 17. This correlation is defined as function of the Re and φ (cRe* $(1+\varphi)^k$), where $10000 \le \text{Re} \le 30000$ and $1\% \le \varphi \le 5\%$.

The third correlation proposed for the calculation of \overline{Nu} along the heated blocks for different spaces between jets D_j is presented as follows:

$$\overline{\text{Nu}} = 0.4471 \text{Re}^{0.7275} (D_i/H)^{-0.3876}$$
 (24)

and it is validated for an error 3% as shown in Figure 18. This correlation is defined as function of the Re and D_j/H (cRe*(D_j/H)*), where $10000 \le \text{Re} \le 30000$ and $\varphi = 5\%$.

The last correlation proposed for the calculation of \overline{Nu} along the heated blocks for different spaces between blocks (D_b) is presented as follows:

$$\overline{\text{Nu}} = 0.1951 \text{Re}^{0.7463} (D_h/H)^{-0.1958}$$
 (25)

and it is validated with the numerical results for an error between -1.5% and 3.5%, as shown in Figure 19. This correlation is defined as function of the Re and D_b/H ($cRe^n(D_b/H)^k$), where $10000 \le Re \le 30000$ and $1cm \le D_b \le 5cm$ for $\varphi = 1\%$.



5. Conclusion

A numerical study revealed the influence of different parameters on the behavior of water- Al_2O_3 nanofluid flow, heat transfer and entropy generation in a symmetrical channel containing heated blocks and cooled by impinging jets. The most important results obtained in the study that we conducted are:

- The increase of Re and φ improves the HTR.
- The decrease of D_b can significantly contribute to increasing the HTR.
- The increase in Re and change the located of the J2 towards the J1 enhanced the HTR, and we got the highest values of HTR when we fixed the position of J2 on the second block.
- The maximum values of $\overline{\mathrm{Nu}}$, located around the last block, is due the domination of the J2 compared to J1.
- TPMM gives higher Nu values than the SPM.
- \bullet Different correlations are proposed for $\overline{\text{Nu}}$.
- The increase in values of Re and φ leads to increase of $\dot{S}_{q,v}$, $\dot{S}_{q,th}$ and $\dot{S}_{q,t}$.
- Be is approximately equal to 1 at the cases, because of the dominance of $\dot{S}_{g,th}$.
- The increase in Re and D_b values increases $\dot{S}_{a,t}$.
- The increase in Re and the change of the located J2 towards the J1 increases the values of \dot{S}_{at} .

Our future perspective based on studying effect of the nanoparticles shape using hybrid nanofluid with different models of turbulence such as SST kw, RSM..., study the effect of some parameters geometric on the HTR such as length and height of hot blocks, we reformulate the current study using more complex 3D geometries.

Author Contributions

The manuscript was written through the contribution of all authors. All authors discussed the results, reviewed, and approved the final version of the manuscript.

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Conflict of Interest

The authors declared no potential conflicts of interest with respect to the research, authorship, and publication of this article.

Nomenclature

Ве	Bejan number	$\dot{S}_{g,t}$	Total entropy generation (W/m³k)
Cf	Friction factor	$\dot{S}_{q,th}$	Thermal entropy generation (W/m³k)
Ср	Specific heat (J/kgk)	$\dot{S}_{q,v}$	Frictional entropy generation (W/m³k)
dp	Nanoparticles diameter (m)	T	Temperature (K)
D_{b}	Spacing between blocks (m)	Abbreviati	ons
D_j	Distance between J1 and J2	HTR	Heat transfer rate
f drag	Drag function	SPM	Single phase model
g	Gravitational acceleration (m/s²)	SPMM	Single phase mixture model
\overline{h}	Heat transfer coefficient (W/m²k)	TPMM	Two phase mixture model
h'	Height block (m)	Greek sym	abols
Н	Channel height (m)	φ	Volume fraction of nanoparticles (%)
k	Thermal conductivity (W/mk)	ho	Density (kg/m³)
L	Length blocks (m)	μ	Dynamic viscosity (kg/ms)
N	Width of the jets inlet (m)	v	Kinematics viscosity (m²/s)
Nu	Nusselt number	Subscripts	
$\overline{\text{Nu}}$	Average Nusselt number	bf	base fluid
P	Pressure (Pa)	m	mixture



Pr Prandtl number np nanoparticle

Re Reynolds number dr,k drift of water or nanoparticle

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